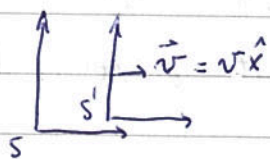


Eletrodinâmica Relativística



$$\vec{\nabla}' \times \vec{E}' = -\frac{\partial \vec{B}'}{\partial t'}, \quad \vec{\nabla}' \times \vec{H}' = \frac{\partial \vec{D}'}{\partial t'} + \vec{j}'$$

$$\vec{\nabla}' \cdot \vec{D}' = \rho', \quad \vec{\nabla}' \cdot \vec{B}' = 0$$

(motuação de índices repetidos) $(\vec{\nabla}' \times \vec{E}')_i = \epsilon_{ijk} \frac{\partial E'_k}{\partial x'_j} = \epsilon_{ijk} \partial'_j E'_k$

$$\vec{\nabla}' \cdot \vec{B}' = \frac{\partial B'_i}{\partial x'_i} = \partial'_i B'_i$$

Soma implícita, $\delta_{ij} x_j = x_i$, $\epsilon_{ijk} \epsilon_{klm} = \delta_{il} \delta_{jm} - \delta_{im} \delta_{jl}$

$$(\vec{\nabla}' \times \vec{E}')_i = \epsilon_{ijk} \partial'_j E'_k = -\partial'_i B'_i, \quad \vec{\nabla}' \cdot \vec{B}' = \partial'_i B'_i = 0$$

seja $F(x', y', z', t')$: $dF = \frac{\partial F}{\partial x'} dx' + \frac{\partial F}{\partial y'} dy' + \frac{\partial F}{\partial z'} dz' + \frac{\partial F}{\partial t'} dt' = \partial'_i F dx'_i$

para um dado evento (x', y', z', t') :

$$dx' = \frac{\partial x'}{\partial x} dx + \frac{\partial x'}{\partial y} dy + \frac{\partial x'}{\partial z} dz + \frac{\partial x'}{\partial t} dt, \text{ etc.}$$

de acordo com as transf. de Lorentz:

$$x' = \gamma(x - vt), \quad y' = y, \quad z' = z$$

$$t' = \gamma(t - vx/c^2), \quad \gamma = (1 - v^2/c^2)^{-1/2} = \text{const.}$$

$$\Rightarrow \frac{\partial x'}{\partial x} = \gamma, \quad \frac{\partial x'}{\partial y} = \frac{\partial x'}{\partial z} = 0, \quad \frac{\partial x'}{\partial t} = -\gamma v \Rightarrow dx' = \gamma dx - \gamma v dt$$

análogamente: $dy' = dy, \quad dz' = dz$

$$dt' = \frac{\partial t'}{\partial x} dx + \frac{\partial t'}{\partial y} dy + \frac{\partial t'}{\partial z} dz + \frac{\partial t'}{\partial t} dt = -\frac{\gamma v}{c^2} dx + \gamma dt$$

Substituindo em dF:

$$dF = \frac{\partial F}{\partial x'} (\gamma dx - \gamma v dt) + \frac{\partial F}{\partial y'} dy + \frac{\partial F}{\partial z'} dz + \frac{\partial F}{\partial t'} (\gamma dt - \frac{\gamma v}{c^2} dx) =$$

$$= (\gamma \frac{\partial F}{\partial x'} - \frac{\gamma v}{c^2} \frac{\partial F}{\partial t'}) dx + \frac{\partial F}{\partial y'} dy + \frac{\partial F}{\partial z'} dz + (\gamma \frac{\partial F}{\partial t'} - \gamma v \frac{\partial F}{\partial x'}) dt \Rightarrow$$

$$\frac{\partial}{\partial x} = \gamma \left(\frac{\partial}{\partial x'} - \frac{v}{c^2} \frac{\partial}{\partial t'} \right), \quad \frac{\partial}{\partial y'} = \frac{\partial}{\partial y}, \quad \frac{\partial}{\partial z'} = \frac{\partial}{\partial z}, \quad \frac{\partial}{\partial t} = \gamma \left(\frac{\partial}{\partial t'} - v \frac{\partial}{\partial x'} \right)$$

$$\partial_x = \gamma (\partial_{x'} - v \partial_{t'}) \quad , \quad \partial_{y'} = \partial_y \quad , \quad \partial_{z'} = \partial_z \quad , \quad \partial_t = \gamma (\partial_{t'} - v \partial_{x'})$$

agora, por exemplo: $\frac{\partial E_x}{\partial z} - \frac{\partial E_z}{\partial x} = -\frac{\partial B_y}{\partial t} \Rightarrow$
(coord. y)

$$\Rightarrow \frac{\partial E_x}{\partial z'} - \gamma \left(\frac{\partial E_z}{\partial x'} - \frac{v}{c^2} \frac{\partial E_z}{\partial t'} \right) = -\gamma \left(\frac{\partial B_y}{\partial t'} - v \frac{\partial B_y}{\partial x'} \right) \Rightarrow$$

$$\Rightarrow \frac{\partial E_x}{\partial z'} - \frac{\partial}{\partial x'} \gamma (E_z + v B_y) = -\frac{\partial}{\partial t'} \gamma (B_y + \frac{v}{c^2} E_z)$$

em S' : $\frac{\partial E_x'}{\partial z'} - \frac{\partial E_z'}{\partial x'} = -\frac{\partial B_y'}{\partial t'} \Rightarrow$

transformações:
$$\begin{cases} E_x' = E_x \\ E_z' = \gamma (E_z + v B_y) \\ B_y' = \gamma (B_y + \frac{v}{c^2} E_z) \end{cases}$$

agora, da coord. z: $\frac{\partial E_y}{\partial x} - \frac{\partial E_x}{\partial y} = -\frac{\partial B_z}{\partial t} \Rightarrow$

$$\Rightarrow \gamma \left(\frac{\partial E_y}{\partial x'} - \frac{v}{c^2} \frac{\partial E_y}{\partial t'} \right) - \frac{\partial E_x}{\partial y'} = -\gamma \left(\frac{\partial B_z}{\partial t'} - v \frac{\partial B_z}{\partial x'} \right) \Rightarrow$$

$$\Rightarrow \frac{\partial}{\partial x'} \gamma (E_y - v B_z) - \frac{\partial E_x}{\partial y'} = -\frac{\partial}{\partial t'} \gamma (B_z - \frac{v}{c^2} E_y)$$

em S' : $\frac{\partial E_y'}{\partial x'} - \frac{\partial E_x'}{\partial y'} = -\frac{\partial B_z'}{\partial t'} \Rightarrow$

transformações:
$$\begin{cases} E_y' = \gamma (E_y - v B_z) \\ E_x' = E_x \\ B_z' = \gamma (B_z - \frac{v}{c^2} E_y) \end{cases}$$

e da coord. x: $\frac{\partial E_z}{\partial y} - \frac{\partial E_y}{\partial z} = -\frac{\partial B_x}{\partial t} \Rightarrow$

$$\Rightarrow \frac{\partial E_z}{\partial y'} - \frac{\partial E_y}{\partial z'} = -\gamma \left(\frac{\partial B_x}{\partial t'} - v \frac{\partial B_x}{\partial x'} \right) \Rightarrow \text{transf. inversas em } \vec{E}_z \text{ e } \vec{E}_y$$

$$\Rightarrow \gamma \left(\frac{\partial E_z'}{\partial y'} - v \frac{\partial B_y'}{\partial z'} \right) - \gamma \left(\frac{\partial E_y'}{\partial z'} + v \frac{\partial B_z'}{\partial y'} \right) = -\gamma \left(\frac{\partial B_x}{\partial t'} - v \frac{\partial B_x}{\partial x'} \right) \Rightarrow$$

$$\Rightarrow -v \left(\frac{\partial B_x}{\partial x'} + \frac{\partial B_y'}{\partial y'} + \frac{\partial B_z'}{\partial z'} \right) + \frac{\partial E_z'}{\partial y'} - \frac{\partial E_y'}{\partial z'} = -\frac{\partial B_x}{\partial t'} \Rightarrow$$

$$\Rightarrow \frac{\partial B_x}{\partial x'} + \frac{\partial B_y'}{\partial y'} + \frac{\partial B_z'}{\partial z'} = \frac{1}{v} \left(\frac{\partial E_z'}{\partial y'} - \frac{\partial E_y'}{\partial z'} + \frac{\partial B_x}{\partial t'} \right) \quad (1)$$

mas de $\vec{\nabla} \cdot \vec{B} = 0 \Rightarrow$

$$\frac{\partial B_x}{\partial x} + \frac{\partial B_y}{\partial y} + \frac{\partial B_z}{\partial z} = 0 \Rightarrow$$

$$\gamma \left(\frac{\partial B_x}{\partial x'} - \frac{v}{c^2} \frac{\partial B_x}{\partial t'} \right) + \frac{B_y}{\partial y'} + \frac{\partial B_z}{\partial z'} = 0 \Rightarrow \text{transf. inversas em } B_y \text{ e } B_z:$$

$$\gamma \left(\frac{\partial B_x}{\partial x'} - \frac{v}{c^2} \frac{\partial B_x}{\partial t'} \right) + \gamma \left(\frac{\partial B_y'}{\partial y'} - \frac{v}{c^2} \frac{\partial E_z'}{\partial y'} \right) + \gamma \left(\frac{\partial B_z'}{\partial z'} + \frac{v}{c^2} \frac{\partial E_y'}{\partial z'} \right) = 0 \Rightarrow$$

$$\frac{\partial B_x}{\partial x'} + \frac{\partial B_y'}{\partial y'} + \frac{\partial B_z'}{\partial z'} = \frac{v}{c^2} \left(\frac{\partial E_z'}{\partial y'} - \frac{\partial E_y'}{\partial z'} + \frac{\partial B_x}{\partial t'} \right) \quad (2)$$

em S' : $\frac{\partial B_x'}{\partial x'} + \frac{\partial B_y'}{\partial y'} + \frac{\partial B_z'}{\partial z'} = 0$, $\frac{\partial E_z'}{\partial y'} - \frac{\partial E_y'}{\partial z'} = -\frac{\partial B_x'}{\partial t'}$

(1) e (2) $\Rightarrow B_x' = B_x$

completando as transformações:

$$\begin{aligned} E_x' &= E_x \\ E_y' &= \gamma(E_y - v B_z) \\ E_z' &= \gamma(E_z + v B_y) \\ B_x' &= B_x \\ B_y' &= \gamma(B_y + \frac{v}{c^2} E_z) \\ B_z' &= \gamma(B_z - \frac{v}{c^2} E_y) \end{aligned}$$

$$\begin{aligned} E_x &= E_x' \\ E_y &= \gamma(E_y' + v B_z') \\ E_z &= \gamma(E_z' - v B_y') \\ B_x &= B_x' \\ B_y &= \gamma(B_y' - \frac{v}{c^2} E_z') \\ B_z &= \gamma(B_z' + \frac{v}{c^2} E_y') \end{aligned}$$

def. do tensor (de rank-2, antissimétrico)
intensidade de campo :

$$F^{\mu\nu} = \begin{pmatrix} 0 & E_x/c & E_y/c & E_z/c \\ -E_x/c & 0 & B_z & -B_y \\ -E_y/c & -B_z & 0 & B_x \\ -E_z/c & B_y & -B_x & 0 \end{pmatrix}$$

o do tensor dual $G^{\mu\nu} = \begin{pmatrix} 0 & B_x & B_y & B_z \\ -B_x & 0 & -E_z/c & E_y/c \\ -B_y & E_z/c & 0 & -E_x/c \\ -B_z & -E_y/c & E_x/c & 0 \end{pmatrix}$

as equações de Maxwell tornam-se:

1) $\partial_\nu G^{\mu\nu} = \frac{\partial G^{\mu\nu}}{\partial x^\nu} = 0$, 2) $\partial_\nu F^{\mu\nu} = \frac{\partial F^{\mu\nu}}{\partial x^\nu} = \mu_0 J^\mu$

1) as equações homogêneas:

$\partial_\nu G^{\mu\nu} = \frac{\partial G^{\mu\nu}}{\partial x^\nu} = 0^\mu = (0, \vec{0}) = 0$
(4 componentes)

para $\mu=0$: $\frac{\partial G^{0\nu}}{\partial x^\nu} = \frac{\partial G^{00}}{\partial x^0} + \frac{\partial G^{01}}{\partial x^1} + \frac{\partial G^{02}}{\partial x^2} + \frac{\partial G^{03}}{\partial x^3} =$

$= \frac{\partial 0}{\partial t} + \frac{\partial B_x}{\partial x} + \frac{\partial B_y}{\partial y} + \frac{\partial B_z}{\partial z} = \boxed{\vec{\nabla} \cdot \vec{B} = 0}$

para $\mu=1$: $\frac{\partial G^{1\nu}}{\partial x^\nu} = \frac{\partial G^{10}}{\partial x^0} + \frac{\partial G^{11}}{\partial x^1} + \frac{\partial G^{12}}{\partial x^2} + \frac{\partial G^{13}}{\partial x^3} =$

$= -\frac{\partial B_x}{\partial t} + 0 - \frac{\partial E_z}{\partial y} + \frac{\partial E_y}{\partial z} = -\frac{1}{c} \left(\frac{\partial \vec{B}}{\partial t} + \vec{\nabla} \times \vec{E} \right)_x = 0$

analogamente: $\mu=2$: $\frac{\partial G^{2\nu}}{\partial x^\nu} = -\frac{1}{c} \left(\frac{\partial \vec{B}}{\partial t} + \vec{\nabla} \times \vec{E} \right)_y = 0$

$\mu=3$: $\frac{\partial G^{3\nu}}{\partial x^\nu} = -\frac{1}{c} \left(\frac{\partial \vec{B}}{\partial t} + \vec{\nabla} \times \vec{E} \right)_z = 0$

$\Rightarrow \boxed{\vec{\nabla} \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}}$

2) equações não-homogêneas

def.: quadrivetor densidade de corrente $J^\mu = (c\rho, \vec{J})$

lembrando que $\vec{J} = \rho \vec{v}$ e que $V = \gamma(c, \vec{v})$

agora $\frac{\partial J^\mu}{\partial x^\mu} = 0$ é a equação da continuidade:

centro $\partial - \frac{\partial J^\mu}{\partial x^\mu}$

$\frac{\partial J^\mu}{\partial x^\mu} = \frac{\partial J^0}{\partial x^0} + \frac{\partial J^1}{\partial x^1} + \frac{\partial J^2}{\partial x^2} + \frac{\partial J^3}{\partial x^3} = \frac{c \partial \rho}{\partial t} + \vec{\nabla} \cdot \vec{J} = 0 \Rightarrow \boxed{\vec{\nabla} \cdot \vec{J} + \frac{\partial \rho}{\partial t} = 0}$

$$\partial_\nu F^{\mu\nu} = \frac{\partial F^{\mu\nu}}{\partial x^\nu} = \mu_0 J^\mu$$

para $\mu=0$:
$$\frac{\partial F^{0\nu}}{\partial x^\nu} = \frac{\partial F^{00}}{\partial x^0} + \frac{\partial F^{01}}{\partial x^1} + \frac{\partial F^{02}}{\partial x^2} + \frac{\partial F^{03}}{\partial x^3} =$$

$$= 0 + \frac{1}{c} \left(\frac{\partial E_x}{\partial x} + \frac{\partial E_y}{\partial y} + \frac{\partial E_z}{\partial z} \right) = \frac{1}{c} \vec{\nabla} \cdot \vec{E}$$

$\mu_0 J^0 = \mu_0 c \rho$

$$\Rightarrow \vec{\nabla} \cdot \vec{E} = \mu_0 c^2 \rho = \frac{\mu_0}{\epsilon_0} \rho = \frac{\rho}{\epsilon_0} \Rightarrow \boxed{\vec{\nabla} \cdot \vec{E} = \rho / \epsilon_0}$$

para $\mu=1$:
$$\frac{\partial F^{1\nu}}{\partial x^\nu} = \frac{\partial F^{10}}{\partial x^0} + \frac{\partial F^{11}}{\partial x^1} + \frac{\partial F^{12}}{\partial x^2} + \frac{\partial F^{13}}{\partial x^3} =$$

$$= -\frac{\partial E_x}{c^2 \partial t} + 0 + \frac{\partial B_z}{\partial y} - \frac{\partial B_y}{\partial z} = \left(-\frac{1}{c^2} \frac{\partial \vec{E}}{\partial t} + \vec{\nabla} \times \vec{B} \right)_x$$

analogamente para $\mu=2,3$:
$$\frac{\partial F^{2\nu}}{\partial x^\nu} = \left(-\frac{1}{c^2} \frac{\partial \vec{E}}{\partial t} + \vec{\nabla} \times \vec{B} \right)_y, \quad \frac{\partial F^{3\nu}}{\partial x^\nu} = \left(-\frac{1}{c^2} \frac{\partial \vec{E}}{\partial t} + \vec{\nabla} \times \vec{B} \right)_z$$

para $\mu=1,2,3$: $\mu_0 J^\mu = \mu_0 J_x, \mu_0 J_y, \mu_0 J_z$

$$\Rightarrow -\frac{1}{c^2} \frac{\partial \vec{E}}{\partial t} + \vec{\nabla} \times \vec{B} = \mu_0 \vec{J} \Rightarrow \boxed{\vec{\nabla} \times \vec{B} = \mu_0 \vec{J} + \frac{1}{c^2} \frac{\partial \vec{E}}{\partial t}}$$

$$\Rightarrow \boxed{\partial_\nu G^{\mu\nu} = 0}, \quad \boxed{\partial_\nu F^{\mu\nu} = \mu_0 J^\mu}$$

As equações de Maxwell ficam:

$$1) \vec{\nabla}' \cdot \vec{E}' = \frac{\rho'}{\epsilon_0}, \quad \begin{aligned} \rho' &= \gamma(\rho - v J_x / c) \Rightarrow \rho' = \gamma(\rho - v J_x \mu_0 \epsilon_0) \\ E_x' &= E_x \\ E_y' &= \gamma(E_y - v B_z) \\ E_z' &= \gamma(E_z + v B_y) \end{aligned} \quad \left\{ \begin{aligned} \partial_x' &= \gamma(\partial_x + \frac{v}{c^2} \partial_t) \\ \partial_y' &= \partial_y, \partial_z' = \partial_z \end{aligned} \right.$$

$$\begin{aligned} \vec{\nabla}' \cdot \vec{E}' &= \partial_x' E_x' + \partial_y' E_y' + \partial_z' E_z' = \gamma(\partial_x E_x + \frac{v}{c^2} \partial_t E_x) + \\ &\quad + \gamma(\partial_y E_y - v \partial_y B_z) + \gamma(\partial_z E_z + v \partial_z B_y) = \\ &= \gamma(\vec{\nabla} \cdot \vec{E}) + \gamma v \left(\frac{1}{c^2} \partial_t E_x - \partial_y B_z + \partial_z B_y \right) = \\ &= \gamma \frac{\rho}{\epsilon_0} + \gamma v \underbrace{(-\mu_0 J_x)}_{(\vec{\nabla} \times \vec{B})_x} = \\ &= \gamma \frac{\rho}{\epsilon_0} - \gamma v J_x \mu_0 \frac{\epsilon_0}{\epsilon_0} = \frac{\gamma(\rho - v J_x \mu_0 \epsilon_0)}{\epsilon_0} = \frac{\rho'}{\epsilon_0} \end{aligned}$$

$$\begin{aligned} 2) \vec{\nabla}' \cdot \vec{B}' &= \partial_x' B_x' + \partial_y' B_y' + \partial_z' B_z' = \gamma(\partial_x B_x + \frac{v}{c^2} \partial_t B_x) + \\ &\quad + \gamma(\partial_y B_y + \frac{v}{c^2} \partial_y E_z) + \gamma(\partial_z B_z - \frac{v}{c^2} \partial_z E_y) = \\ &= \gamma \vec{\nabla} \cdot \vec{B} + \frac{\gamma v}{c^2} (\partial_t B_x + \partial_y E_z - \partial_z E_y) \stackrel{0}{=} 0 \\ &\quad \underbrace{(\vec{\nabla} \times \vec{E})_x}_{= -\partial_t B_x} \end{aligned}$$

$$\begin{aligned} 3) (\vec{\nabla}' \times \vec{E}')_x &= \partial_y' E_z' - \partial_z' E_y' = \\ &= \gamma(\partial_y E_z + v \partial_y B_y) - \gamma(\partial_z E_y - v \partial_z B_z) = \\ &= \gamma(\partial_y E_z - \partial_z E_y) + \gamma v (\partial_y B_y + \partial_z B_z) = \\ &= \gamma (\vec{\nabla} \times \vec{E})_x + \gamma v (-\partial_x B_x) = -\gamma \partial_t B_x - \gamma v \partial_x B_x \end{aligned}$$

$$\text{mas } \partial_t' B_x' = \gamma(\partial_t + v \partial_x) B_x = \gamma \partial_t B_x + \gamma v \partial_x B_x$$

$$\therefore (\vec{\nabla}' \times \vec{E}')_x = -\partial_t' B_x'$$

$$A) \vec{\nabla}' \times \vec{B}' = \mu_0 \vec{J}' + \frac{1}{c^2} \frac{\partial \vec{E}'}{\partial t'}$$

$$J_x' = \gamma (J_x - v/c^2 \rho)$$

$$\partial t' = \gamma (\partial t + v \partial x)$$

$$(\vec{\nabla}' \times \vec{B}')_x = \partial_y' B_z' - \partial_z' B_y' =$$

$$= \gamma (\partial_y B_z - v/c^2 \partial_y E_y) - \gamma (\partial_z B_y + v/c^2 \partial_z E_z) =$$

$$= \gamma (\partial_y B_z - \partial_z B_y) - \gamma v/c^2 (\partial_y E_y + \partial_z E_z) =$$

$$= \gamma (\vec{\nabla} \times \vec{B})_x - \gamma v/c^2 (\rho/\epsilon_0 - \partial_x E_x) =$$

$$= \gamma \mu_0 J_x + (\gamma v/c^2) \partial_t E_x + (\gamma v/c^2) \partial_x E_x - \gamma v \rho / c^2 \epsilon_0 =$$

$$= \mu_0 \gamma (J_x - v \rho) + \gamma v/c^2 (\partial_t E_x + \partial_x E_x) =$$

$$= \mu_0 J_x' + \frac{1}{c^2} \partial_t' E_x'$$