

Símbolos de Christoffel a partir das equações de Euler-Lagrange

$$L = g_{\mu\nu} \frac{dx^\mu}{d\lambda} \frac{dx^\nu}{d\lambda} \quad [\text{veja eq. 3.3.14 do Wald}]$$

\* Tomemos como exemplo a métrica  $ds^2 = -\left(1 - \frac{2M}{r}\right) dt^2 + \left(1 + \frac{2M}{r}\right) dr^2 + r^2(d\theta^2 + \sin^2\theta d\phi^2)$

$$\Rightarrow L = -\left(1 - \frac{2M}{r}\right) \left(\frac{dt}{d\lambda}\right)^2 + \left(1 + \frac{2M}{r}\right) \left(\frac{dr}{d\lambda}\right)^2 + r^2 \left(\frac{d\theta}{d\lambda}\right)^2 + r^2 \sin^2\theta \left(\frac{d\phi}{d\lambda}\right)^2$$

\* Equações de Euler Lagrange:  $\frac{\partial L}{\partial x^\mu(\lambda)} - \frac{d}{d\lambda} \left( \frac{\partial L}{\partial \dot{x}^\mu(\lambda)} \right) = 0$

Por exemplo, para  $\mu=0$  temos  $\frac{\partial L}{\partial t} - \frac{d}{d\lambda} \left( \frac{\partial L}{\partial \dot{t}} \right) = 0$

Como  $L$  não depende explicitamente de  $t$ , temos que  $\frac{\partial L}{\partial t} = 0$ . Por outro lado,

$$\frac{\partial L}{\partial \dot{t}} = \frac{\partial L}{\partial \left(\frac{dt}{d\lambda}\right)} = -\left(1 - \frac{2M}{r}\right) \cdot 2 \cdot \left(\frac{dt}{d\lambda}\right) \Rightarrow \frac{d}{d\lambda} \left( \frac{\partial L}{\partial \dot{t}} \right) = -2 \frac{d}{d\lambda} \left[ \left(1 - \frac{2M}{r}\right) \cdot \frac{dt}{d\lambda} \right]$$

$$= -2 \cdot \frac{2M}{r^2} \frac{dr}{d\lambda} \cdot \frac{dt}{d\lambda} - 2 \cdot \left(1 - \frac{2M}{r}\right) \frac{d^2 t}{d\lambda^2}$$

repetir o mesmo para  $\mu=1, \mu=2$  e  $\mu=3$ .

Daí fica:  $\frac{\partial L}{\partial t} - \frac{d}{d\lambda} \left( \frac{\partial L}{\partial \dot{t}} \right) = 0 \Rightarrow -2 \cdot \frac{2M}{r^2} \cdot \frac{dr}{d\lambda} \cdot \frac{dt}{d\lambda} - 2 \cdot \left(1 - \frac{2M}{r}\right) \cdot \frac{d^2 t}{d\lambda^2} = 0$

$$\Rightarrow \frac{d^2 t}{d\lambda^2} + \frac{\frac{2M}{r^2}}{1 - \frac{2M}{r}} \cdot \frac{dr}{d\lambda} \cdot \frac{dt}{d\lambda} = 0 \Rightarrow \boxed{\frac{d^2 t}{d\lambda^2} + \frac{2M}{r^2 - 2Mr} \cdot \frac{dr}{d\lambda} \cdot \frac{dt}{d\lambda} = 0}$$

essa é a componente 0 da equação geodésica  $\frac{d^2 x^\alpha}{d\lambda^2} + \Gamma^\alpha_{\mu\nu} \frac{dx^\mu}{d\lambda} \frac{dx^\nu}{d\lambda} = 0$

Comparando com  $\frac{d^2 x^\alpha}{d\lambda^2} + \Gamma^\alpha_{\mu\nu} \frac{dx^\mu}{d\lambda} \frac{dx^\nu}{d\lambda} = 0 = \frac{d^2 t}{d\lambda^2} + \Gamma^0_{00} \left(\frac{dt}{d\lambda}\right)^2 + \Gamma^0_{11} \left(\frac{dr}{d\lambda}\right)^2 + \Gamma^0_{22} \left(\frac{d\theta}{d\lambda}\right)^2 + \Gamma^0_{33} \left(\frac{d\phi}{d\lambda}\right)^2 + 2\Gamma^0_{01} \frac{dt}{d\lambda} \frac{dr}{d\lambda}$

+  $2\Gamma^0_{02} \frac{dt}{d\lambda} \frac{d\theta}{d\lambda} + 2\Gamma^0_{03} \frac{dt}{d\lambda} \frac{d\phi}{d\lambda} + 2\Gamma^0_{12} \frac{dr}{d\lambda} \frac{d\theta}{d\lambda} + 2\Gamma^0_{13} \frac{dr}{d\lambda} \frac{d\phi}{d\lambda} + 2\Gamma^0_{23} \frac{d\theta}{d\lambda} \frac{d\phi}{d\lambda} \Rightarrow \Gamma^0_{00} = \Gamma^0_{11} = \Gamma^0_{22} = \Gamma^0_{33} = 0$  e  $\Gamma^0_{01} = \frac{M}{r^2 - 2Mr}$